

## On the constitutive assumptions for the recombination term for the R-D-D equations for scintillators

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**Abstract:** The evolution equations for inorganic scintillators (which converts ionizing radiations into visible light) were obtained in [1] (see also [2], [3]): they are Reaction-Diffusion-Drift equations, coupled with the Poisson equation of electrostatics and Neumann boundary conditions, in terms of the  $m$ -dimensional *charge carriers vector*  $n \equiv (n_1, \dots, n_m)$ :

$$\begin{aligned} \operatorname{div}(D[\nabla n] + MN[q \otimes \nabla \varphi]) + r(n) &= \dot{n}, \\ -\epsilon \Delta \varphi &= eq \cdot n, \end{aligned} \quad (1)$$

here  $D$  and  $M$  are the  $m \times m$  *diffusivity* and *mobility* semi-positive definite matrices,  $N(n) = \operatorname{diag}(n_1, \dots, n_m)$ ,  $\varphi$  is the *electric potential*,  $\epsilon$  the *permittivity* of the crystal,  $e$  the elementary charge,  $q \in \mathbb{Z}^m$  the *charge numbers vector* and finally  $r(n)$  the recombination term which accounts for the visible photons production. Equation (1)<sub>1</sub> was obtained by the means of a microstructured continuum model, in the sense of [4], where the rate-of-change of the  $m$ -dimensional *director* is the *scintillation potential*  $g(n) = eq\varphi(n) + F(n)$ , where  $F(n) > 0$  is a term of entropic nature.

Here we shall deal with various constitutive assumptions for the recombination term  $r(n)$  and show how, when  $F(n)$  is the Gibbs-Boltzmann entropy, we can recover the cubic polynomial expression used *e.g.* in one of the few previous existing phenomenological models [5]. Further we shall show how a polynomial representation can be obtained as a limit of a Markov process and also explore the instance in which  $F(n)$  is represented instead by the Fermi-Dirac potentials.

### References

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